

# *Beyond Chern-Simons: Historical & Background Experimental techniques Challenges*

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11/5/24



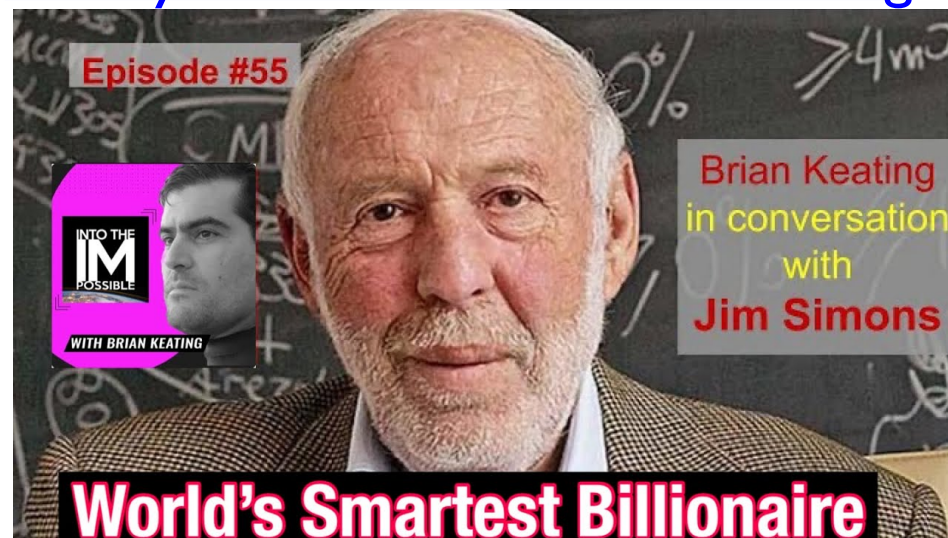
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# Outline



Cosmology with CMB Polarization



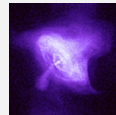
Cosmic birefringence



Experimental Results

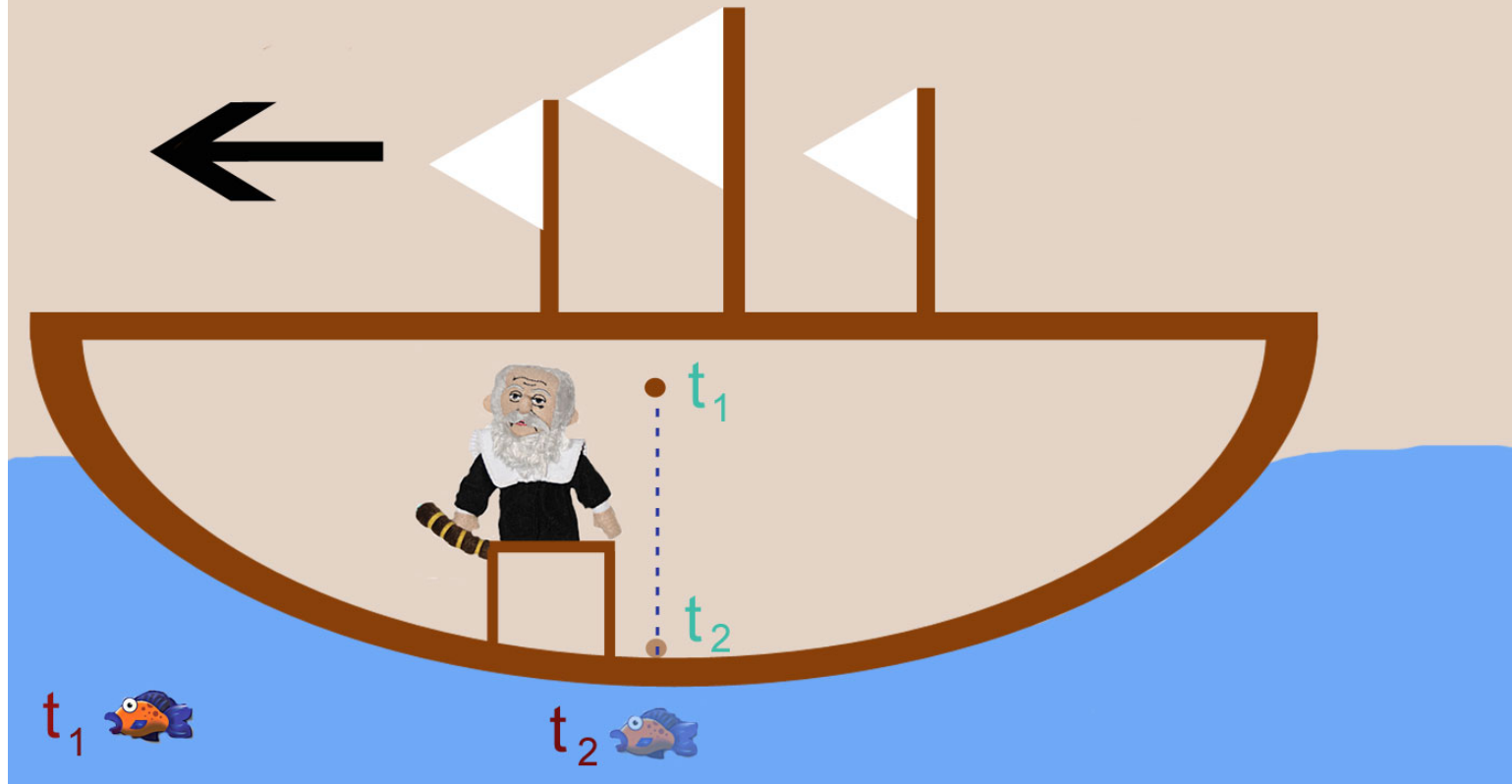


PolCal Drones!



Future

# Galileo Galilei: Moral Relativist

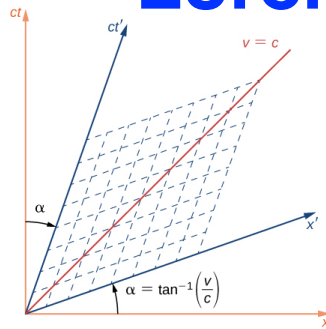


Galileo used this experiment to demonstrate that Earth could, in fact, move without conflicting with our everyday sense of the world. Even if the Earth were moving, everything on its surface would move at the same speed. Consequently, doing these experiments in the geocentric universe of the Church and the heliocentric universe Galileo had in mind would be indistinguishable. Galileo had stumbled upon a type of relativity that now bears his own name. Reflecting upon several simple mechanical experiments like the falling ball, Galileo concluded: You will discover not the least change in all the effects named, nor could you tell from any of them whether the ship was moving or standing still.

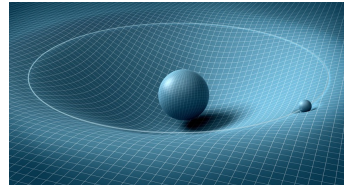
<https://www.physicscentral.com/explore/plus/galilean-relativity.cfm>



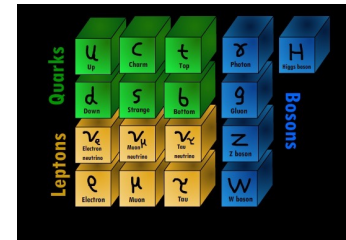
# Lorentz invariance violation?



General Relativity (GR)



Standard Model (SM)

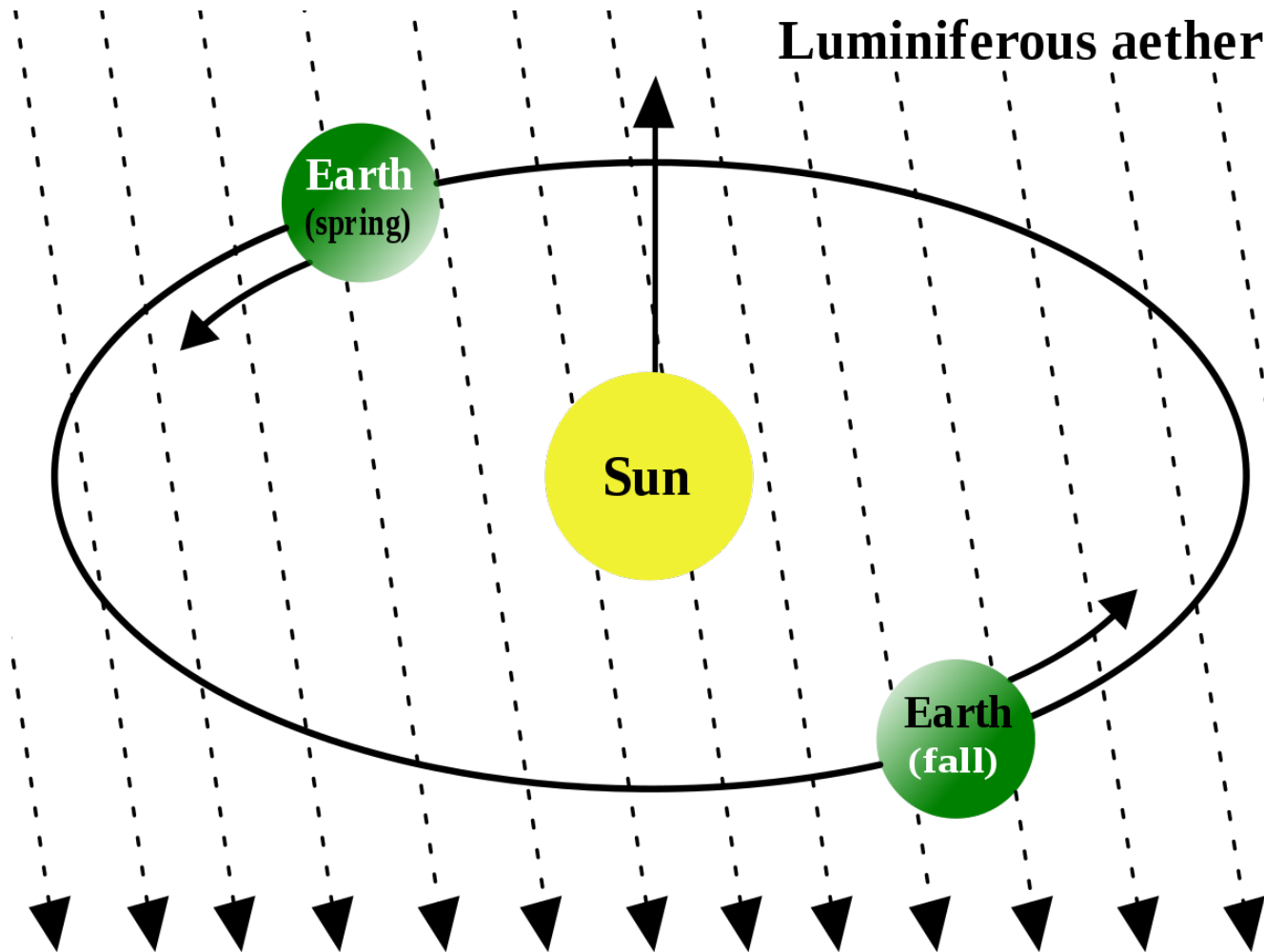


- Both GR and SM assume Lorentz Symmetry (special relativity, relativistic quantum theory, QFT)
- But GR and Quantum Field Theory/Standard Model are not unified!

**Violations of Lorentz symmetry, if detected, would point us toward new physics, quantum gravity!**

# Michelson-Morley: Hold my aether!

**Luminiferous aether**





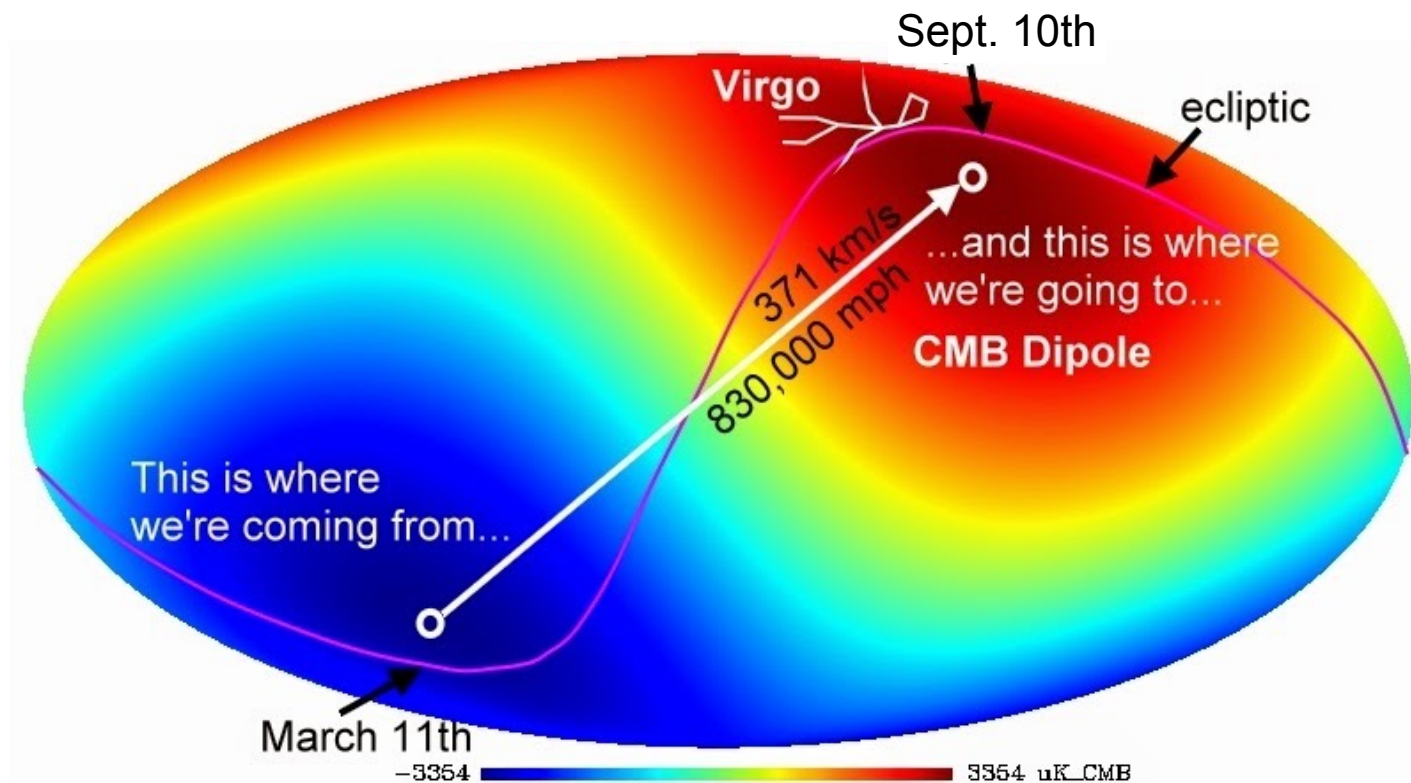


Are we sure there isn't a preferred reference frame?



# CMB DIPOLE

- **CMB temp. map has dipole distortion.**
- Each point on sky has blackbody spectrum, but spectrum is redshifted/blueshifted in different halves of the sky.
- Due to Doppler shift caused by net motion of satellite with respect to frame in which CMB is isotropic. **Net motion toward Virgo cluster!**





# Is Parity violated? **Yes!**

- Cobalt-60 cooled to 3 mK (!)
- Applied magnetic field to align all atoms
- Measured electron emission
- Reversed the magnetic field, reversing the spin
- Saw the same anisotropic electron emission
- Preferred emission direction is opposite to the nuclear spin
- Thus we can tell a “mirrored” world from our own – parity NOT conserved

PHYSICAL REVIEW

VOLUME 104, NUMBER 1

OCTOBER 1, 1956

COSMO-LOGICAL



The question of p  
experiments are sug

ained. Possible

## Experimental Test of Parity Conservation in Beta Decay\*

C. S. Wu, *Columbia University, New York, New York*

AND

E. AMBLER, R. W. HAYWARD, D. D. HOPPE, AND R. P. HUDSON,  
*National Bureau of Standards, Washington, D. C.*

(Received January 15, 1957)

IN a recent paper<sup>1</sup> on the question of parity in weak interactions, Lee and Yang critically surveyed the experimental information concerning this question and reached the conclusion that there is no existing evidence either to support or to refute parity conservation in weak interactions. They proposed a number of experiments on beta decays and hyperon and meson decays which would provide the necessary evidence for parity conservation or nonconservation. In beta decay, one could measure the angular distribution of the electrons coming from beta decays of polarized nuclei. If an asymmetry in the distribution between  $\theta$  and  $180^\circ - \theta$  (where  $\theta$  is the angle between the orientation of the parent nuclei and the momentum of the electrons) is observed, it provides unequivocal proof that parity is not conserved in beta decay. This asymmetry effect has been observed in the case of oriented  $\text{Co}^{60}$ .

It has been known for some time that  $\text{Co}^{60}$  nuclei can be polarized by the Rose-Gorter method in cerium magnesium (cobalt) nitrate, and the degree of polarization detected by measuring the anisotropy of the succeeding gamma rays.<sup>2</sup> To apply this technique to the present problem, two major difficulties had to be over-

sotropy alone provides a reliable measure of nuclear polarization. Specimens were made by taking good single crystals of cerium magnesium nitrate and growing on the upper surface only an additional crystalline layer containing  $\text{Co}^{60}$ . One might point out here that since the allowed beta decay of  $\text{Co}^{60}$  involves a change of spin of

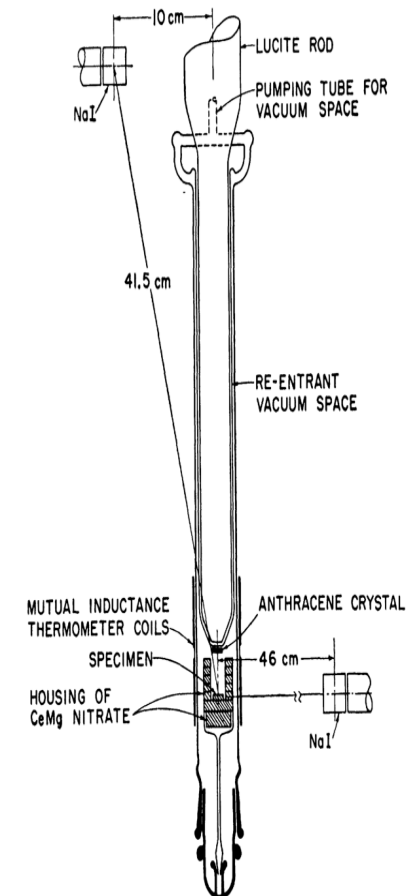


FIG. 1. Schematic drawing of the lower part of the cryostat.

Carroll, Field & Jackiw (1990); Harari & Sikivie (1992); Carroll (1998)

# Cosmic Birefringence

The effect accumulates over the distance



- If the Universe is filled with a pseudo-scalar field (e.g., an axion field) coupled to the electromagnetic tensor via a Chern-Simons coupling:

**Turner & Widrow (1988)**

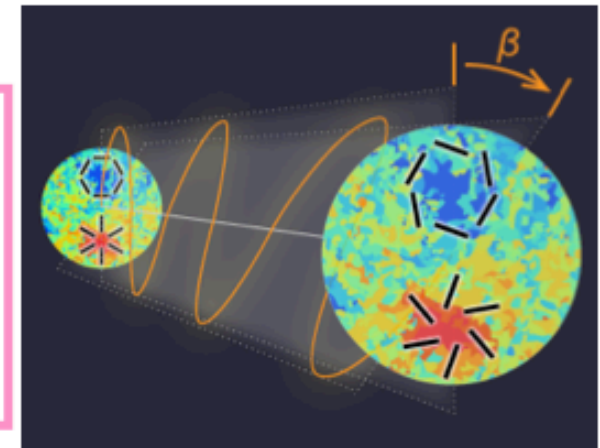
the effective Lagrangian for axion electrodynamics is

$$\mathcal{L} = -\frac{1}{2}\partial_\mu\theta\partial^\mu\theta - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \underbrace{g_a\theta F_{\mu\nu}\tilde{F}^{\mu\nu}}_{\text{Chern-Simons term}}, \quad (3.7)$$

$\tilde{F}^{\mu\nu} = \sum_{\alpha\beta} \frac{\epsilon^{\mu\nu\alpha\beta}}{2\sqrt{-g}} F_{\alpha\beta}$

where  $g_a$  is a coupling constant of the order  $\alpha$ , and the vacuum angle  $\theta = \phi_a / f_a$  ( $\phi_a$  = axion field). The equations

$$\beta = 2g_a \int_{t_{\text{emission}}}^{t_{\text{observed}}} dt \dot{\theta}$$



The larger the distance the photon travels, the larger the effect becomes.



**HAPPY BIRTHDAY TO CARROLL, FIELD, & JACKIW (1990)**  
**(Galileo Galilei's 426th BIRTHDAY 🎉)**

PHYSICAL REVIEW D

VOLUME 41, NUMBER 4

15 FEBRUARY 1990

**Limits on a Lorentz- and parity-violating modification of electrodynamics**

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(Received 5 September 1989)

The Chern-Simons Lagrangian has been studied previously in  $(2+1)$ -dimensional spacetime, where it is both gauge and Lorentz invariant. In  $3+1$  dimensions, this term couples the dual electromagnetic tensor to an external four-vector. If we take this four-vector to be fixed, the term is gauge invariant but not Lorentz invariant. In this paper, we examine both the theoretical consequences of such a modification and observational limits we can put on its magnitude. The Chern-Simons term would rotate the plane of polarization of radiation from distant galaxies, an effect which is not observed. From the observations we deduce that the magnitude of the vector is  $< 1.7 \times 10^{-42} h_0$  GeV, where  $h_0$  is the Hubble constant in units of  $100 \text{ km sec}^{-1} \text{ Mpc}^{-1}$ .



<https://www.researchgate.net/publication/13277928> Limits on a Lorentz- and parity-violating modification of electrodynamics

## I. INTRODUCTION

Gauge and Lorentz invariance are two symmetries of Maxwell's electrodynamics that have come to dominate all fundamental physical theory. They provide physical principles that guide the invention of models describing fundamental phenomena, and their experimental status—within electromagnetism—is well established. The properties of electromagnetic radiation, both in a natural setting and in high-energy accelerators, are precisely described by Lorentz-invariant dynamics. Gauge invariance, interpreted as the masslessness of the photon, is validated by stringent limits on the photon mass.

Experimental tests of such well-established and universal physical ideas are best discussed within a theoretical framework that allows departures to be governed by arbitrary parameters; experimental data then set limits on the magnitude of these symmetry-breaking parameters. Thus, violations of gauge invariance are parametrized by a mass  $\mu$  for the photon field  $A_\nu$ . A mass term is hypothesized to modify the electromagnetic Maxwell Lagrange density  $\mathcal{L}_{\text{EM}}$ ,

$$\mathcal{L}_{\text{EM}} = -\frac{1}{4} F_{\nu\lambda} F^{\nu\lambda} \quad (1)$$

so that the photon becomes massive:

$$\mathcal{L}_\mu = -\frac{1}{4} F_{\nu\lambda} F^{\nu\lambda} + \frac{\mu^2}{2} A^\nu A_\nu. \quad (2)$$

Here  $F_{\nu\lambda}$  is the electromagnetic tensor  $F_{\nu\lambda} = \partial_\nu A_\lambda - \partial_\lambda A_\nu$ , and the field equations in the presence of a conserved current  $J_\nu$  read

$$\square A_\nu + \mu^2 A_\nu = 4\pi J_\nu, \quad \partial_\nu A^\nu = 0, \quad (3)$$

where  $\square$  is the d'Alembertian  $\square = \partial_t^2 - \nabla^2$ . (We set  $c$  equal to unity throughout.) Gauge invariance

$$A_\nu \rightarrow A_\nu + \partial_\nu \chi \quad (4)$$

is clearly lost. Geomagnetic data then set the limit<sup>1</sup>

$\mu \leq 3 \times 10^{-24}$  GeV; observations of the galactic magnetic field set the more stringent bound<sup>2</sup> of  $\mu \leq 3 \times 10^{-36}$  GeV; see below.

In this paper we explore the experimental limits on another modification of Maxwell theory, which also involves a mass parameter  $p_\alpha$ , but respects gauge invariance—rather, it is Lorentz invariance that is violated.

The modification we consider involves adding to the Maxwell Lagrange density a Chern-Simons term:

$$\mathcal{L}_p = \mathcal{L}_{\text{EM}} + \mathcal{L}_{\text{CS}}. \quad (5)$$

The Chern-Simons term is given by

$$\mathcal{L}_{\text{CS}} = -\frac{1}{2} p_\alpha A_\beta \tilde{F}^{\alpha\beta}, \quad (6)$$

where  $\tilde{F}^{\alpha\beta}$  is the dual electromagnetic tensor,  $\tilde{F}^{\alpha\beta} = \frac{1}{2} \epsilon^{\alpha\beta\mu\nu} F_{\mu\nu}$ . This modification couples the electromagnetic field to an (as yet unspecified) four-vector  $p_\alpha$ .

When electromagnetic phenomena are confined to a plane, as in the quantum Hall effect and high- $T_c$  superconductivity, the approximation can be made that no interesting dynamical motion takes place in the direction perpendicular to the plane. Then the external vector  $p_\alpha$  may be chosen to lie in that direction as well, and (6) reduces to an unconventional electrodynamic action that is Lorentz and gauge invariant in a three-dimensional spacetime, i.e., boosts in the plane leave dynamics unchanged. It was in this context that the Chern-Simons term was initially investigated as a “topological mass” term for gauge fields in (2+1)-dimensional spacetime.<sup>3</sup> Models in which  $\mathcal{L}_{\text{CS}}$  is taken to be the entire gauge field action have found application in examinations of the quantum Hall effect<sup>4</sup> and high- $T_c$  superconductivity.<sup>5</sup> Moreover, several purely mathematical applications for  $\mathcal{L}_{\text{CS}}$  have also been found.<sup>6</sup>

In this paper we shall consider the (3+1)-dimensional case, where considerations of both Lorentz and gauge invariance play a crucial role.

James Clerk Maxwell

(born June 13, 1831, Edinburgh, Scotland—died November 5, 1879, Cambridge, Cambridgeshire, England)

The **action**  $S$  of Chern–Simons theory is proportional to the integral of the **Chern–Simons 3-form**

$$S = \frac{k}{4\pi} \int_M \text{tr} (A \wedge dA + \frac{2}{3} A \wedge A \wedge A).$$

The constant  $k$  is called the *level* of the theory. The classical physics of Chern–Simons theory is independent of the choice of level  $k$ .

Classically the system is characterized by its equations of motion which are the extrema of the action with respect to variations of the field  $A$ . In terms of the field curvature

$$F = dA + A \wedge A$$

the **field equation** is explicitly

$$0 = \frac{\delta S}{\delta A} = \frac{k}{2\pi} F.$$

where  $K$  is the Chern-Simons three-form:

**Can break both parity and Lorentz Invariance symmetry**

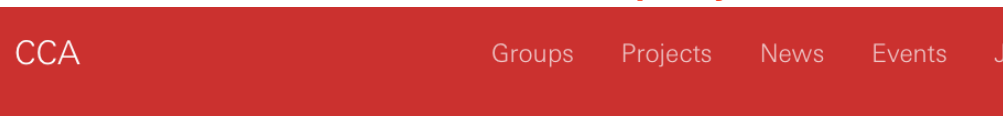
CS Action

$$S_{CS} = p \times F \wedge F$$

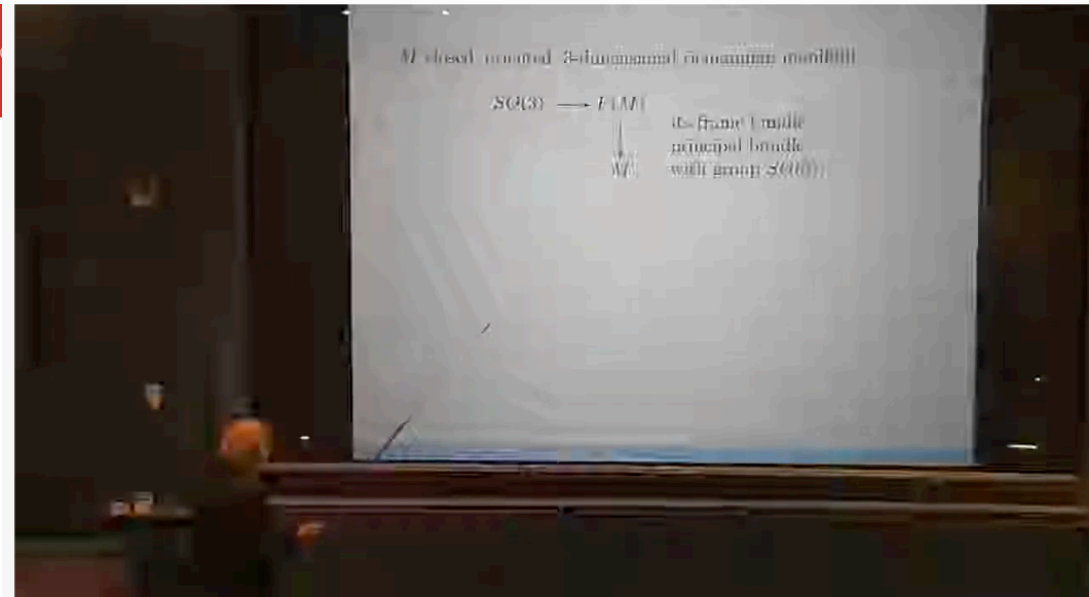
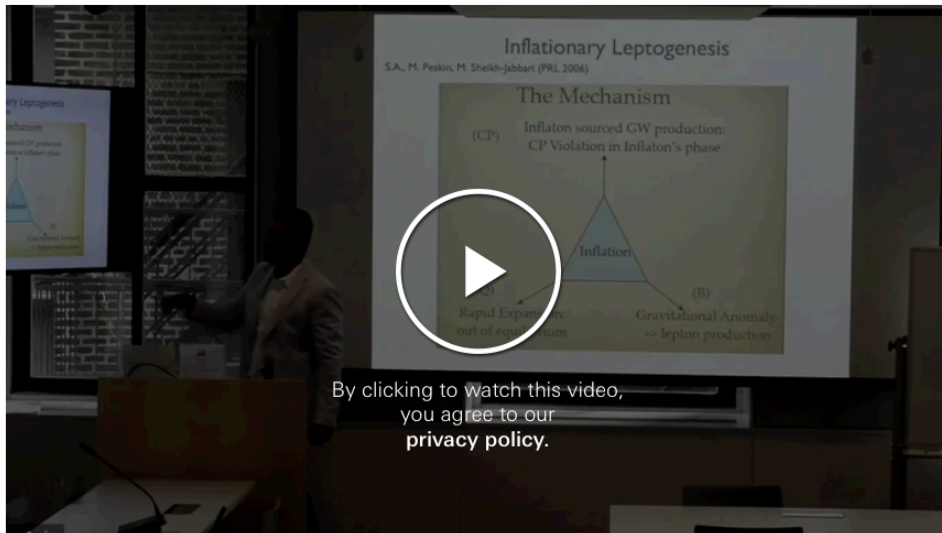
$$\vec{p} = p_\mu = (p_0, \vec{p})$$

$$F \wedge F = dK$$

$$K = A \wedge F$$



## Stephon Alexander: The Theory of a Halo: Chern Simons Theory – Dark Genesis and Inflation



C. Vafa, M. Freedman, C. Kane, Jim Simons - Applications of Chern-Simons Theory (April 25, 2018)

906 views • Feb 6, 2019

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# 1. Introduction and Preliminaries

## Electromagnetic Field Tensor and Forms

- ▶ Electromagnetic 2-form:  $F = \frac{1}{2} F_{\mu\nu} dx^\mu \wedge dx^\nu$ , with  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$
- ▶ Hodge dual:  $\tilde{F} = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} F_{\alpha\beta} dx_\mu \wedge dx_\nu$
- ▶ 4D Action Term:

$$S = \int d^4x P(x) F \wedge F \tag{1}$$

- ▶  $P(x)$  is a pseudoscalar field,  $F = dA$



## 2. Chern-Simons Current and 3D Reduction

### Rewriting the Action

- ▶  $F \wedge F = dK$ , where  $K = A \wedge F$
- ▶  $S = \int d^4x dP \wedge K$
- ▶ For  $P(t)$  (time-dependent):

$$S = \int d^4x (\partial_0 P(t)) \epsilon^{ijk} A_i F_{jk}$$

### Physical Implication

- ▶ Modification to photon propagation in 4D space-time
- ▶ Not symmetric under spatial parity, leading to polarization asymmetry

### 3. Physical Interpretation and Final Form

#### **Electromagnetic Interpretation**

- ▶  $K^0 = \epsilon^{ijk} A_i F_{jk}$  as magnetic helicity density
- ▶ Topological measure of magnetic field linkage and twist

#### **Axion-Coupled Action**

- ▶  $S = - \int d^4x (\partial_0 a(t)) K^0$
- ▶ Observable effect: cosmic birefringence (rotation of CMB polarization)

# Birefringence Generating Mechanisms

New fields (such as axions) with Chern-Simons couplings

Polarization rotation for a single photon:

- $\Delta\theta \propto g_{\phi\gamma}(\phi_{\text{absorbed}} - \phi_{\text{emitted}})$

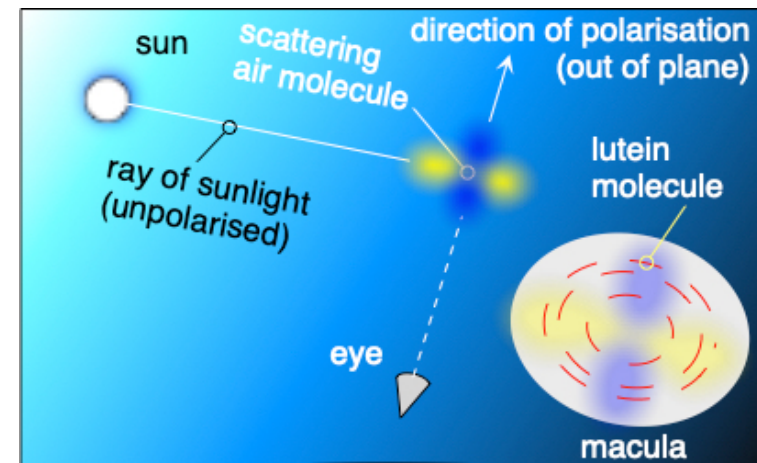
Some axion masses cause polarization oscillation on the timescale of hours - months

particle (e.g. axion) field      EM field

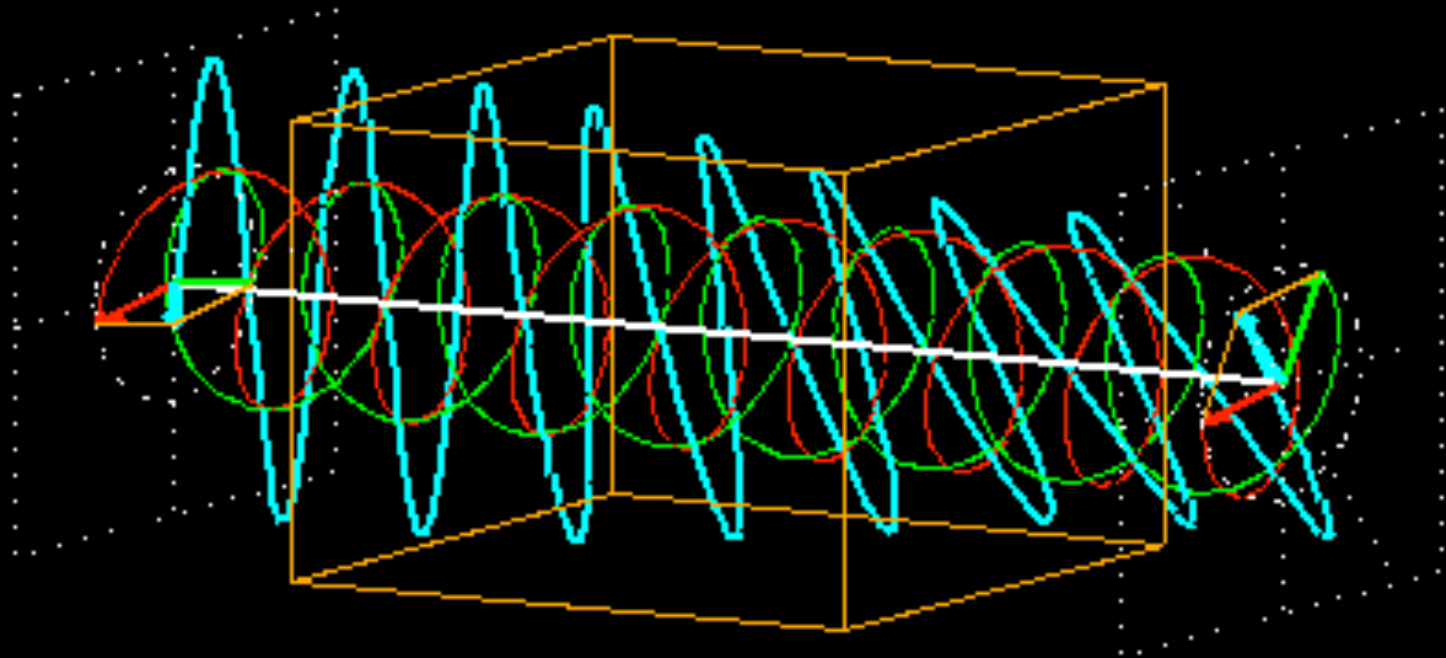
coupling

$$-\frac{1}{4}g_{\phi\gamma}\phi F_{\mu\nu}\tilde{F}^{\mu\nu}$$

Chern-Simons term in the action  
coupling EM field and axion field



# Can we test Lorentz Invariance with the CMB?



**Look for a preferred direction in spacetime!**  
*(Discovered already???)*





# 'This Side Up' May Apply To the Universe, After All

Continued From Page A1

ward Sextans or Aquila, would be a matter of arbitrary choice.

The discovery was made by Dr. Borge Nodland of Rochester and Dr. John Ralston of Kansas, using radio-wave observations made by different astronomers around the world. In a report to be published on Monday in Physical Review Letters, the two physicists concluded on a note of excitement tempered with caution.

"Barring hidden systematic bias in the data," they wrote, the behavior of electromagnetic radiation propagating over vast distances "indicates a new cosmological effect."

In an announcement by the University of Rochester yesterday, Dr. Nodland said: "The big news is that perhaps not all space is equal, for as far back as we can peer in time. This work defies the notion that there is no 'up' or 'down' in space."

Dr. Ralston said, "Our observational data suggest that there is a mysterious axis, a kind of cosmological north star that orients the universe."

Few other physicists and cosmologists have had a chance to read the journal report, but they agreed that the research must be tested thoroughly before the conclusions can be accepted.

"It would be a really profound change in physics, if it is true," said Dr. P. James E. Peebles, a Princeton University astrophysicist.

Dr. Stephen P. Maran, an astronomer

at the Goddard Space Flight Center in Greenbelt, Md., said: "Anytime you find a new effect globally in the sky, the crucial issue is always whether you have correctly taken account of systematic errors in the observations. And any result of this potential magnitude is going to be viewed with considerable skepticism until new experiments can be done to verify it."

In their report, Dr. Nodland and Dr. Ralston constructed a mathematical theory that could explain the observations. The data indicate that light actually travels through space at two slightly different speeds. Such a mismatch in speeds would cause the polarization plane to rotate in a certain familiar manner. It is the way physics students see when they pass light through corn syrup and look at the light with polarizing filters.

The physicists say the axis of orientation they have inferred would appear to be along different lines in different parts of the universe, but they would be parallel to the one observed from Earth.

Because the findings run counter to the idea that all space is uniform and that the speed of light in a vacuum is always the same, the implications of the research could be enormous. For example, scientists might have to reconsider the concept that the Big Bang, the theorized moment of cosmic origin, was completely symmetric.

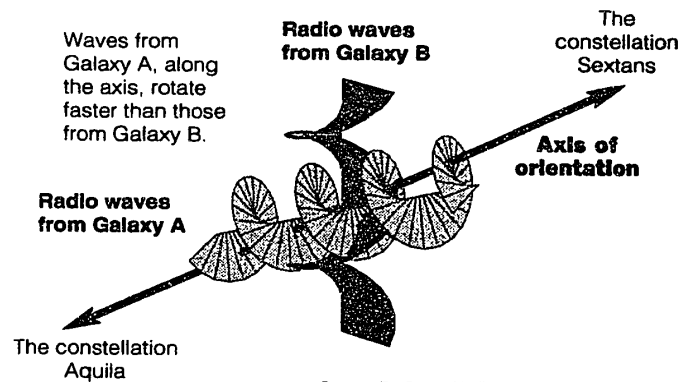
"Perhaps it was not a perfect Big Bang, but a Big Bang with a twist to

## A CLOSER LOOK

### A Universe With Ups and Downs?

Scientists have discovered that radio waves traveling in one direction across the universe behave differently from radio waves traveling at right angles to that direction. This implies that the universe is not uniform in all directions. Astrophysicists have generally assumed it is uniform.

Surprisingly, the radio waves corkscrew faster when they travel along an apparent axis of orientation, running in a direction roughly between the constellation Sextans and the constellation Aquila, with Earth in the middle. This could mean that there is an "up" and a "down" in the universe, but which is which remains purely arbitrary.



Source: Dr. Borge Nodland, University of Rochester

The New York Times

space and time," Dr. Ralston said. "Such a twist would be seen today as a ripple of nonuniformity, perhaps as the axis represents."

Dr. Nodland also speculated that the observed rotations could be the first evidence for physicists who have theorized the existence of other universes. If the universe in which people live was asymmetric at cre-

ation, he said, it raises the possibility of the simultaneous creation of another universe with an opposite twist.

Concluding a summary of the findings for the Internet, Dr. Nodland said, "At this point, the question of what is truly underlying the effect we see is as wide open as space itself."

# Master Equation: cosmic birefringence x 4

$\gamma(\omega, t, \hat{n})$ : Coupling terms

- Cross-terms between different effects
- Example:  $\gamma \propto \boldsymbol{\alpha}(\hat{n}) \cdot \hat{n} \times \omega^n$
- Higher-order corrections

The total rotation angle including all possible birefringence effects can be expressed as:

$$\Delta\phi(\vec{r}, \hat{n}, \omega, t) = \int_0^r \left[ \underbrace{\alpha_0}_{\text{constant isotropic}} + \underbrace{\boldsymbol{\alpha}(\hat{n})}_{\text{direction-dependent}} \cdot \hat{n} + \underbrace{\beta(\omega)}_{\text{frequency-dependent}} + \underbrace{\alpha(t)}_{\text{time-varying}} + \underbrace{\gamma(\omega, t, \hat{n})}_{\text{coupling terms}} \right] dr'$$

$\alpha_0$ : Constant isotropic birefringence

- Uniform rotation independent of direction and frequency
- Simplest case of cosmic birefringence
- Units: [rad/Mpc]

$\boldsymbol{\alpha}(\hat{n})$ : Anisotropic birefringence tensor

- Direction-dependent rotation
- Can be decomposed into spherical harmonics
- $\boldsymbol{\alpha}(\hat{n}) = \sum_{\ell m} a_{\ell m} Y_{\ell m}(\hat{n})$

$\alpha(t)$ : Time-varying component

- Can be periodic:  $\alpha(t) = \alpha_1 \cos(\Omega t)$
- Or monotonic:  $\alpha(t) = \alpha_1 t$

$\beta(\omega)$ : Frequency-dependent term

- Accounts for dispersive effects
- Generally modeled as power law:  $\beta(\omega) \propto \omega^n$
- Important for multi-frequency observations

## Notes

- Integration is along the photon path from emission to observation
- $\hat{n}$  represents the direction of observation
- $\omega$  is the photon frequency
- Effects can be constrained through CMB observations

# Example 1: Constant Isotropic cosmic birefringence

$\gamma(\omega, t, \hat{n})$ : Coupling terms

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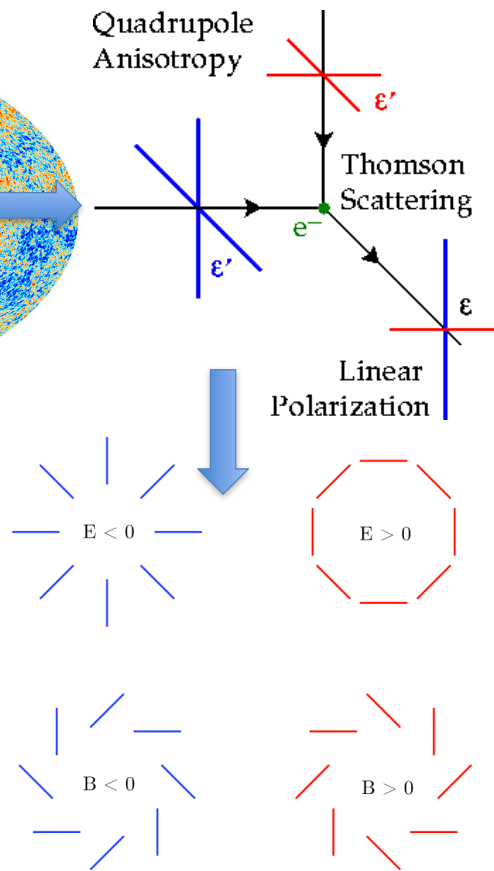
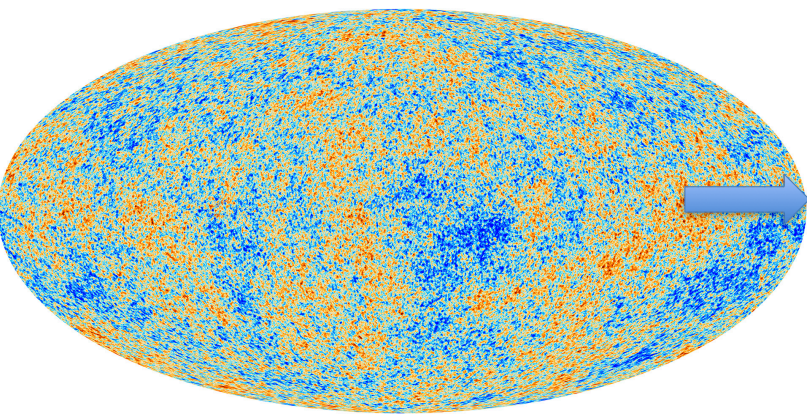
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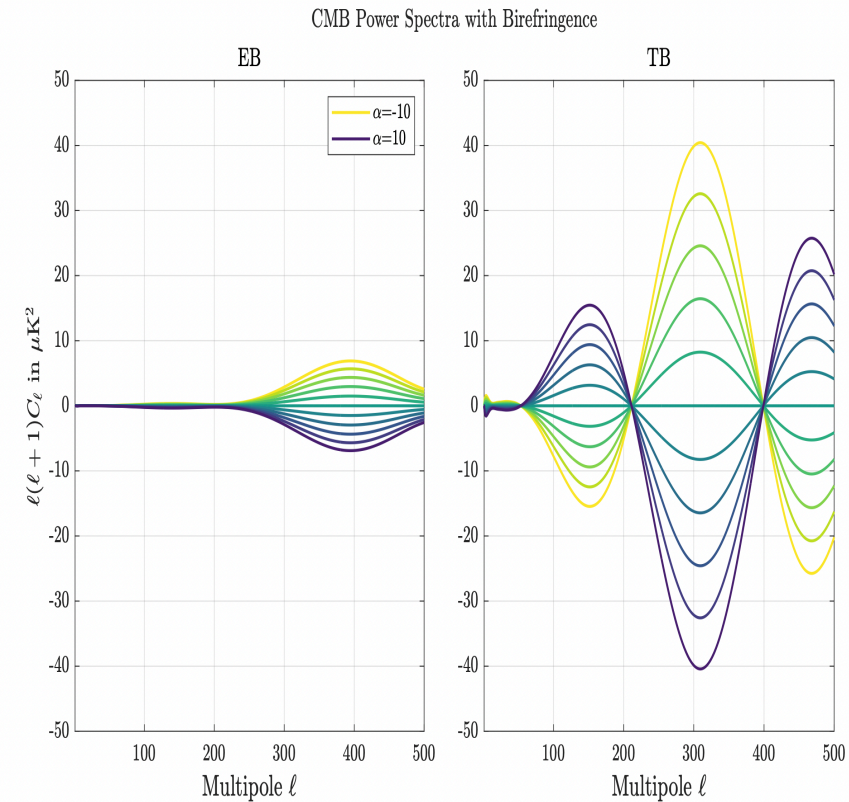
# Cosmic Microwave Background polarization



No TB or EB correlations in the standard model

$$C_{\ell}^{'TB} = C_{\ell}^{TE} \sin(2\alpha)$$

$$C_{\ell}^{'EB} = \frac{1}{2} (C_{\ell}^{EE} - C_{\ell}^{BB}) \sin(4\alpha)$$



**BICEP/Keck XVIII: Measurement of BICEP3 polarization angles and consequences for constraining cosmic birefringence and inflation**  
<https://arxiv.org/pdf/2410.12089>

August 2009

Xia et al. claim a *first* detection of CPT violation!  
Parameterized by Chern-Simons rotation angle  $\alpha$



Contents lists available at ScienceDirect

Physics Letters B

www.elsevier.com/locate/physletb



## Probing CPT violation with CMB polarization measurements

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### ARTICLE INFO

#### Article history:

Received 12 October 2009

Received in revised form 9 March 2010

Accepted 10 March 2010

Available online 16 March 2010

Editor: M. Trodden

#### Keywords:

Observational cosmology

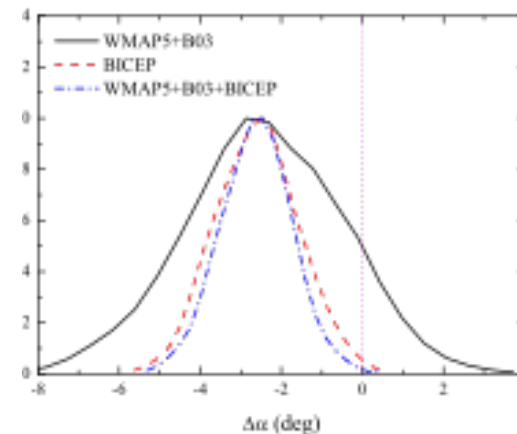
Lorentz and Poincaré invariance

Charge conjugation, parity, time reversal, and other discrete symmetries

### ABSTRACT

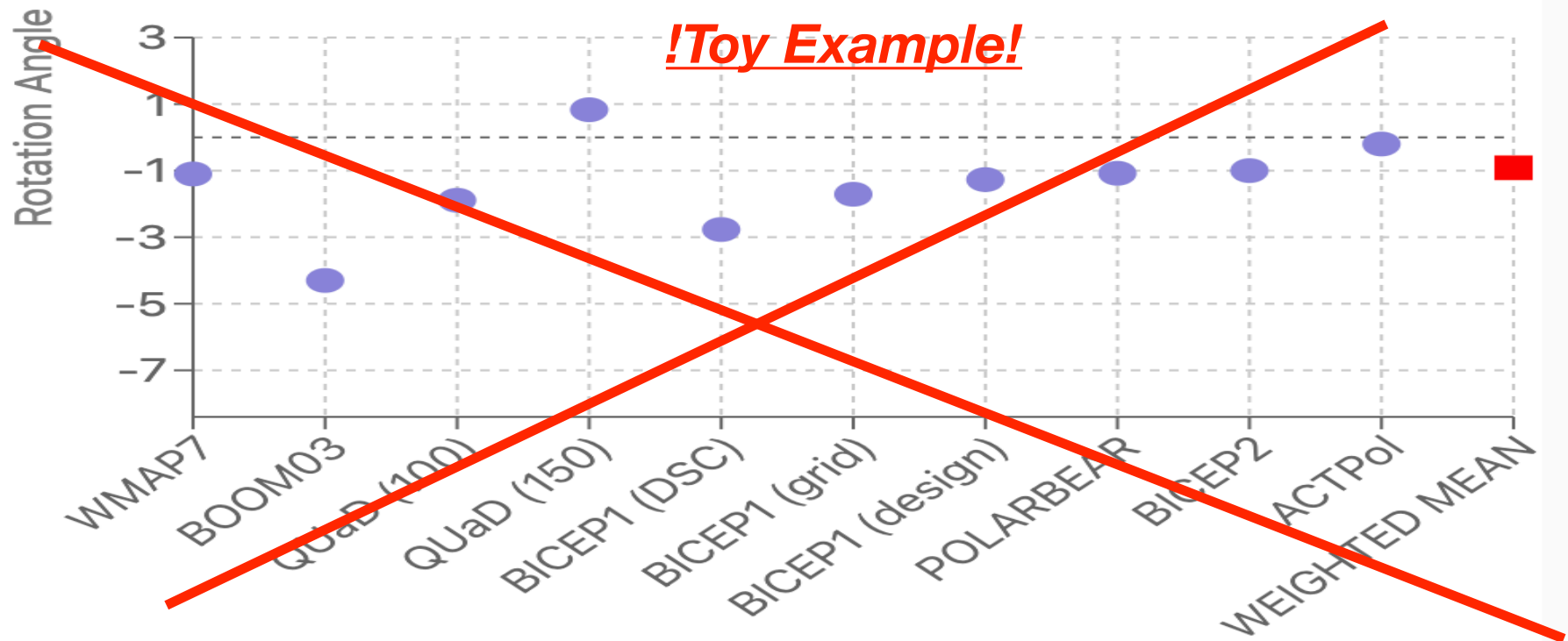
The electrodynamics modified by the Chern-Simons term  $\mathcal{L}_{CS} \sim p_\mu A_\nu \tilde{F}^{\mu\nu}$  with a non-vanishing  $p_\mu$  violates the Charge-Parity-Time Reversal symmetry (CPT) and rotates the linear polarizations of the propagating Cosmic Microwave Background (CMB) photons. In this Letter we measure the rotation angle  $\Delta\alpha$  by performing a global analysis on the current CMB polarization measurements from the five-year Wilkinson Microwave Anisotropy Probe (WMAP5), BOOMERanG 2003 (B03), BICEP and Quid using a Markov Chain Monte Carlo method. Neglecting the systematic errors of these experiments, we find that the results from WMAP5, B03 and BICEP all are consistent and their combination gives  $\Delta\alpha = -2.62 \pm 0.87 \text{ deg}$  (68% C.L.), indicating a  $3\sigma$  detection of the CPT violation. The Quid data alone give  $\Delta\alpha = 0.59 \pm 0.42 \text{ deg}$  (68% C.L.) which has an opposite sign for the central value and smaller error bar compared to that obtained from WMAP5, B03 and BICEP. When combining all the polarization data together, we find  $\Delta\alpha = 0.09 \pm 0.36 \text{ deg}$  (68% C.L.) which significantly improves the previous constraint on  $\Delta\alpha$  and test the validity of the fundamental CPT symmetry at a higher level.

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ne-dimensional posterior distributions of the rotation angle  $\Delta\alpha$  derived from various data combinations. The dotted line illustrates the unrotated case ( $\Delta\alpha = 0$ ) to

## CMB Rotation Angle Measurements from Different Experiments



**Assumes all are independent, ignores systematics etc...**  
**5.4 sigma (BEFORE Minami and Komatsu 2020)**

Error bars show total uncertainty (statistical and systematic combined in quadrature where both are available)

Weighted mean:  $-0.914^\circ \pm 0.169^\circ$  ( $\chi^2/\text{dof} = 7.6/9$ )



**Parity Violation in E&M: Minami & Komatsu PRL 2020:**  
***New extraction of the cosmic birefringence from the Planck 2018 polarization data***

$\beta = +0.35 \pm 0.14$  deg (68%CL)....excludes  $\beta = 0$  @ 99.2%~CL.

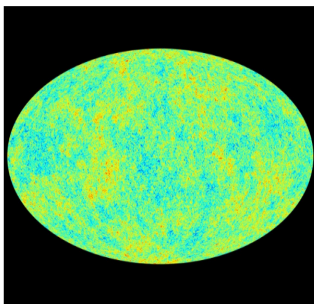
nature > news > article

NEWS • 24 NOVEMBER 2020

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Cosmologists suggest that an exotic substance called quintessence could be accelerating the Universe's expansion – but the evidence is still tentative.

Davide Castelvecchi



A map of the Universe's cosmic microwave background radiation, measured by the Planck space observatory. Credit: ESA

PDF version

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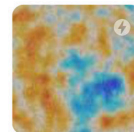
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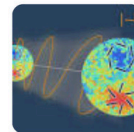
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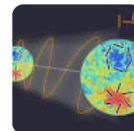
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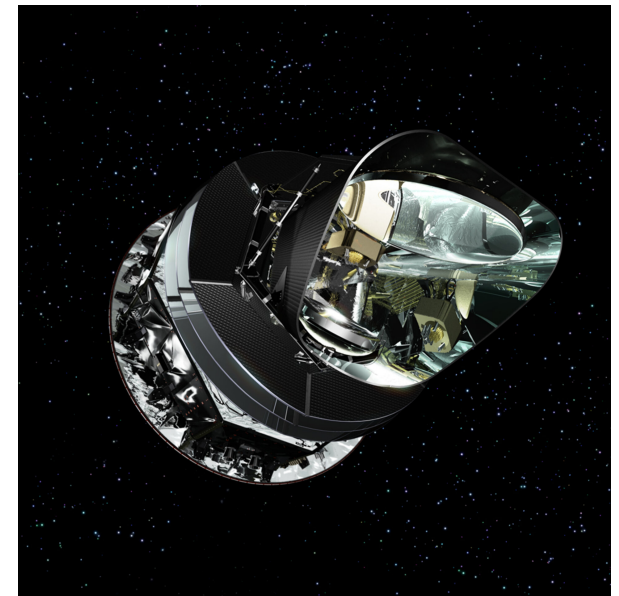
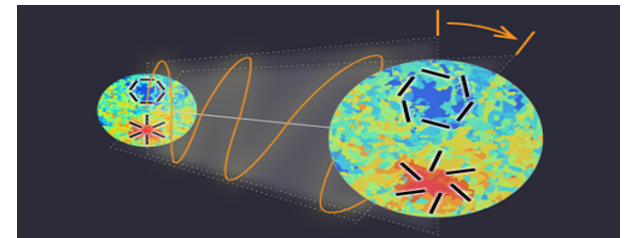
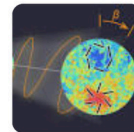
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### A hint of new physics in polarized radiation from the early Universe

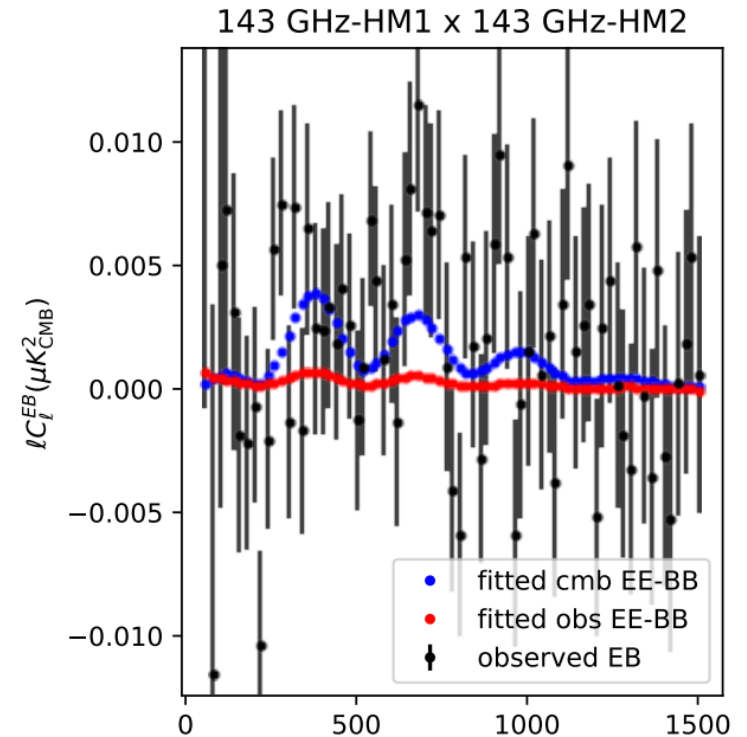
Nov 23, 2020



Planck (ESA)

# Minami & Komatsu (2020) using Planck 2018 data

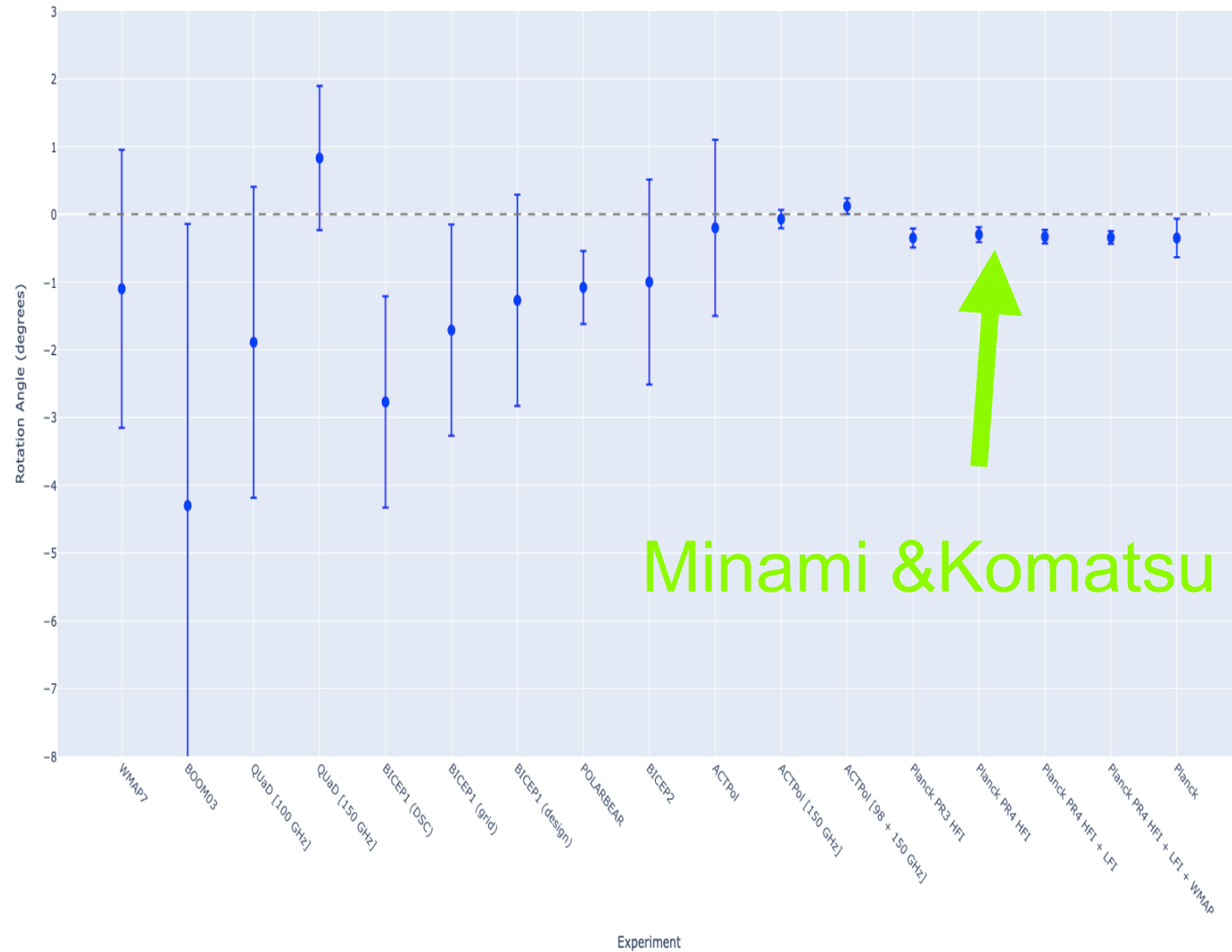
- Isotropic rotation of each CMB photon's polarization plane:
- $0.35^\circ \pm 0.14^\circ$  (**excludes zero at  $2.4\sigma$  level/99.2%**)
- Distinguishes between CMB birefringence and telescope angle miscalibration.
- Key idea: galactic foregrounds are local; should not be rotated by birefringence mechanism.



Part of fig. 2 from *Minami and Komatsu*.  
Blue: correlation with CMB birefringence angle. Red: correlation with telescope miscalibration angle.

Rotation Angle Measurements Across Experiments

(c) @DrBrianKeating 2024



*M&K 2020 Pipeline validation did not include any foregrounds: "Since the FFP10 simulation does not have foreground maps convolved with realistic beam effects such as the  $I \rightarrow P$  leakage, we only consider CMB and noise realizations of the HM maps."*

*Their validation only shows that these foreground-less simulations didn't yield any bias in  $\alpha_v$  or  $\beta$ .*

**This likelihood is a long stretch from the 5-sigma confidence needed for a discovery, but next-generation CMB detectors using the new method could strengthen the result.**

**Data on my website:  
<https://briankeating.com/cb>**

https://briankeating.com/cb

Experiment/Dataset	Frequency [GHz]	$\ell$ range	$\alpha \pm \text{stat}(\pm \text{syst})^{[c]}$	Measurement Method
QUaD[26]	100	200–2000	$-1.89 \pm 2.24(\pm 0.5)$	Polarized source
	150		$+0.83 \pm 0.94(\pm 0.5)$	
BOOM03[27]	143	150–1000	$-4.3 \pm 4.1(\pm 0.69)$	Pre-flight polarized source
ACTPol	146	500–2000	$-0.2 \pm 0.5(-1.2)$	As-designed
WMAP9[28]	23–94	2–800	$0.36 \pm 1.24(\pm 1.5)$	Pre-launch polarized source / Tau A
BICEP2[29]	150	30–300	$-1 \pm 0.2(\pm 1.5)$	Dielectric Sheet
BICEP1[30]	100+150	30–300	$-2.77 \pm 0.86(\pm 1.3)$	Dielectric sheet
			$-1.71 \pm 0.86(\pm 1.3)$	Polarized source
			$-1.08 \pm 0.86(\pm 1.3)$	As-designed
POLARBEAR[31]	150	500–2100	$-1.08 \pm 0.2(\pm 0.5)$	Tau A
Planck[32]	30–353	100–1500	$-0.35 \pm 0.05(\pm 0.28)$	Pre-flight source / Tau A [33, 34]
ACTPol (Choi et al., Murphy et al.)[14, 15]	150	600–1800	$-0.07 \pm 0.09(\pm \sim 0.1)$	Metrology+modeling+point sources
ACTPol (Namikawa et al., Murphy et al. ) [15, 25]	98 + 150	200–2048	$0.12 \pm 0.06(\pm \sim 0.1)$	Metrology+modeling+point sources
Planck PR3 HFI (Minami et al.)[19]	100–353	50–1500	$-0.35 \pm 0.14$	Galactic foregrounds
Planck PR4 HFI (Diego-Palazuelos et al.)[20]	100–353	50–1500	$-0.30 \pm 0.11$	Galactic foregrounds
Planck PR4 HFI + LFI (Eskilt et al.)[21]	30–353	50–1500	$-0.33 \pm 0.10$	Galactic foregrounds
Planck PR4 HFI + LFI + WMAP (Eskilt et al.)[22]	23–353	50–1500	$-0.342^{+0.094}_{-0.091}$	Galactic foregrounds
BICEP3 2-year (this work)	95	40–500	$\alpha \pm 0.078(\pm 0.3)$	Polarized source
Forecast: BICEP3 7-year + RPS improved performance	95	40–500	$\alpha \pm 0.055(\pm \sim 0.07)$	Polarized source

BICEP/Keck XVIII: Measurement of BICEP3 polarization angles and consequences for constraining cosmic birefringence and inflation

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We use a custom-made calibrator to measure individual detectors’ polarization angles of BICEP3, a small aperture telescope observing the cosmic microwave background (CMB) at 96GHz from the South Pole. We describe our calibration strategy and the statistical and systematic uncertainties associated with the measurement. We reach an unprecedented precision for such measurement on a CMB experiment, with a repeatability for each detector pair of 0.02°. We show that the relative angles measured using this method are in excellent agreement with those extracted from CMB data. Because the absolute measurement is currently limited by a systematic uncertainty, we do not derive cosmic birefringence constraints from BICEP3 data in this work. Rather, we forecast the sensitivity of BICEP3 sky maps for such analysis. We investigate the relative contributions of instrument noise, lensing, and dust, as well as astrophysical and instrumental systematics. We also explore the constraining power of different angle estimators, depending on analysis choices. We establish that the BICEP3 2-year dataset (2017–2018) has an on-sky sensitivity to the cosmic birefringence angle  $\alpha_c \approx 0.078^\circ$ , which could be improved to  $\alpha_c \approx 0.055^\circ$  by adding all of the existing BICEP3 data (through 2023). Furthermore, we emphasise the possibility of using the BICEP3 sky patch as a polarization calibration source for CMB experiments, which with the present data could reach a precision of 0.003°. Finally, in the context of inflation searches, we investigate the impact of detector-to-detector variations in polarization angles as they may bias the tensor-to-scalar ratio  $r$ . We show that while the effect is expected to remain subdominant to other sources of systematic uncertainty, it can be reliably calibrated using polarization angle measurements such as the ones we present in this paper.

BICEP3 (James C.)

arXiv:2410.12089v1 [astro-ph.CO] 15 Oct 2024



# Example 2: Time-varying isotropic cosmic birefringence

$\gamma(\omega, t, \hat{n})$ : Coupling terms

- Cross-terms between different effects
- Example:  $\gamma \propto \alpha(\hat{n}) \cdot \hat{n} \times \omega^n$
- Higher-order corrections

The total rotation angle including all possible birefringence effects can be expressed as:

$$\Delta\phi(\vec{r}, \hat{n}, \omega, t) = \int_0^r \left[ \underbrace{\alpha_0}_{\text{constant isotropic}} + \underbrace{\alpha(\hat{n})}_{\text{direction-dependent}} \cdot \hat{n} + \underbrace{\beta(\omega)}_{\text{frequency-dependent}} + \underbrace{\alpha(t)}_{\text{time-varying}} + \underbrace{\gamma(\omega, t, \hat{n})}_{\text{coupling terms}} \right] dr'$$

$\alpha_0$ : Constant isotropic birefringence

- Uniform rotation independent of direction and frequency
- Simplest case of cosmic birefringence
- Units: [rad/Mpc]

$\alpha(\hat{n})$ : Anisotropic birefringence tensor

- Direction-dependent rotation
- Can be decomposed into spherical harmonics
- $\alpha(\hat{n}) = \sum_{\ell m} a_{\ell m} Y_{\ell m}(\hat{n})$

$\alpha(t)$ : Time-varying component

- Can be periodic:  $\alpha(t) = \alpha_1 \cos(\Omega t)$
- Or monotonic:  $\alpha(t) = \alpha_1 t$

$\beta(\omega)$ : Frequency-dependent term

- Accounts for dispersive effects
- Generally modeled as power law:  $\beta(\omega) \propto \omega^n$
- Important for multi-frequency observations

## Notes

- Integration is along the photon path from emission to observation
- $\hat{n}$  represents the direction of observation
- $\omega$  is the photon frequency
- Effects can be constrained through CMB observations

# Axion Signal in the CMB

- Oscillating classical field description:

$$\phi(\vec{x}, t) = \phi_0(\vec{x}, t) \sin(m_\phi t + \theta(\vec{x}))$$

Oscillation period: days-months

- Cosmic birefringence:

$$\beta = \frac{g_{\phi\gamma}}{2} (\phi(\vec{x}_{\text{abs}}, t_{\text{abs}}) - \phi(\vec{x}_{\text{emit}}, t_{\text{emit}}))$$

- CMB polarization angle rotation: (Federreke et. al., PRD 2019)

$$\beta_{\text{CMB}}(t) = \frac{g_{\phi\gamma}\phi_0}{2} \sin(m_\phi t + \theta)$$

Sinusoidal rotation effect: **this work**

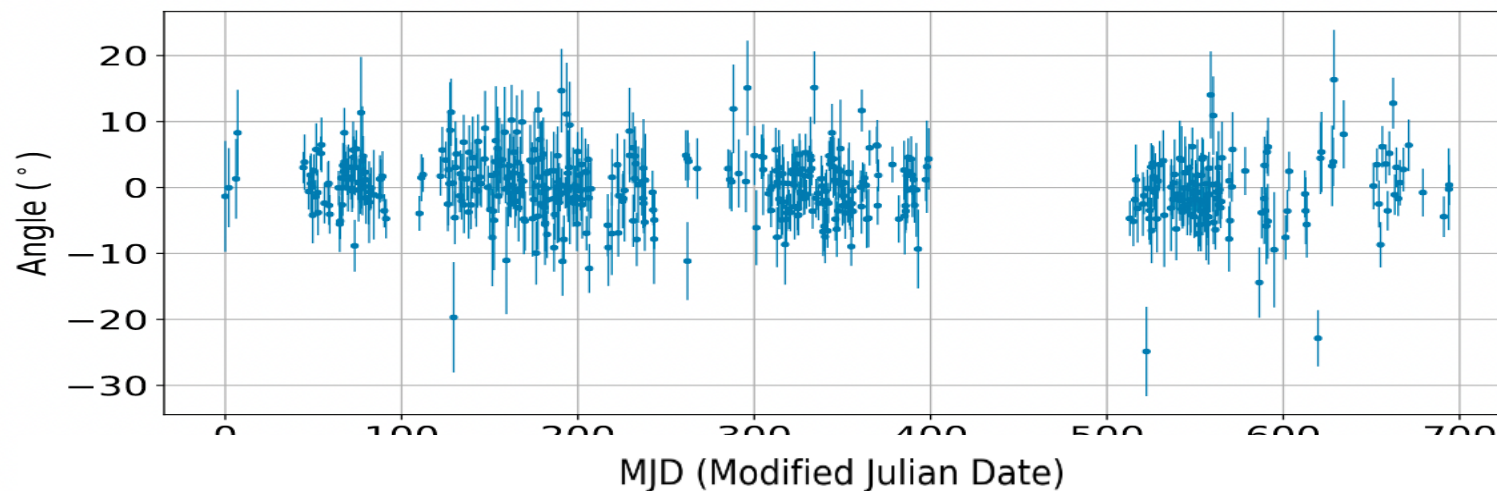
CMB absorbed  
by POLARBEAR 1



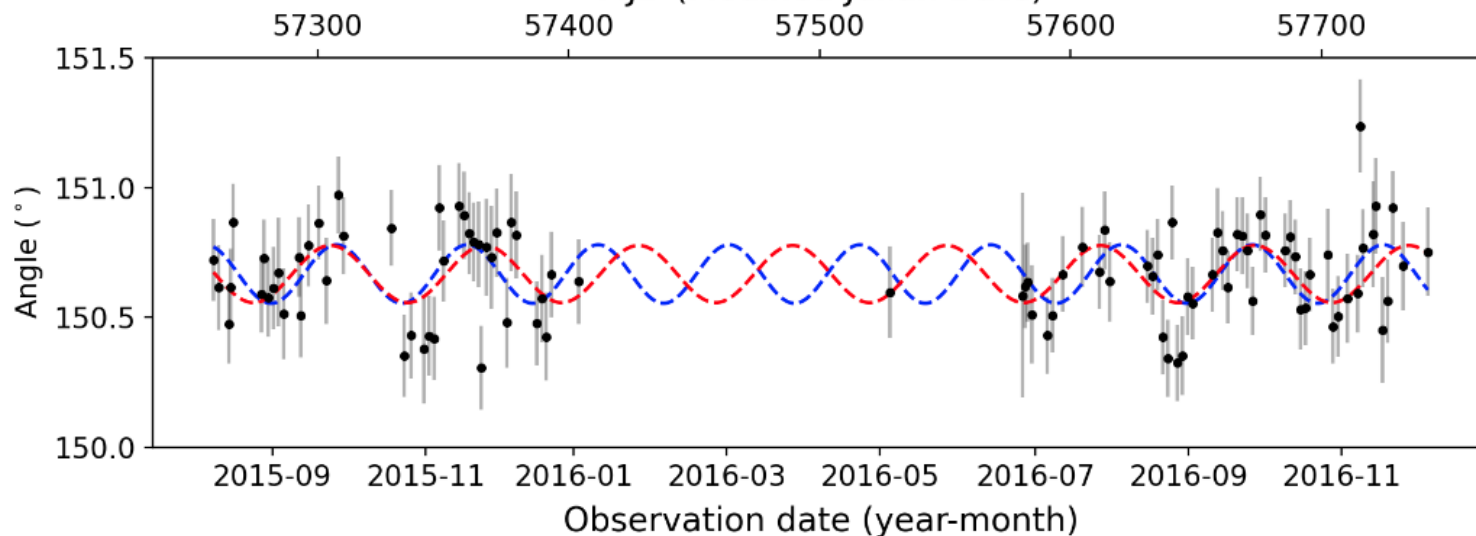
6523 Lyrs



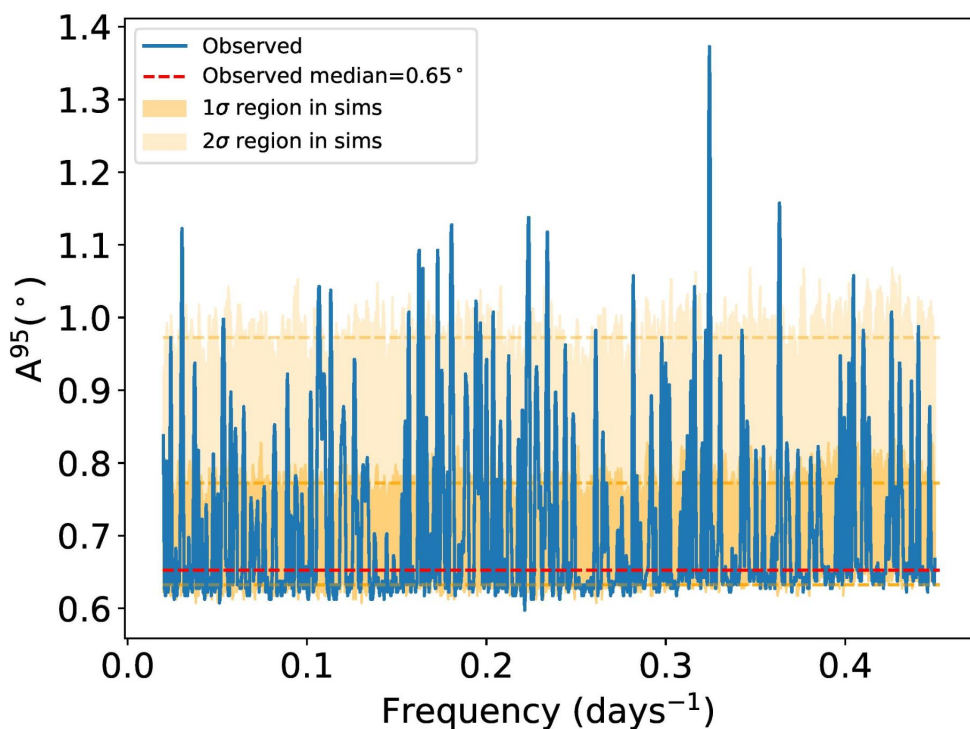
# Estimating an Angle for Each Observation



- Under small rotation angle  $\alpha$ , the maps are:
- Correlate single B map with coadded E maps to estimate angle



# Results: No Detection



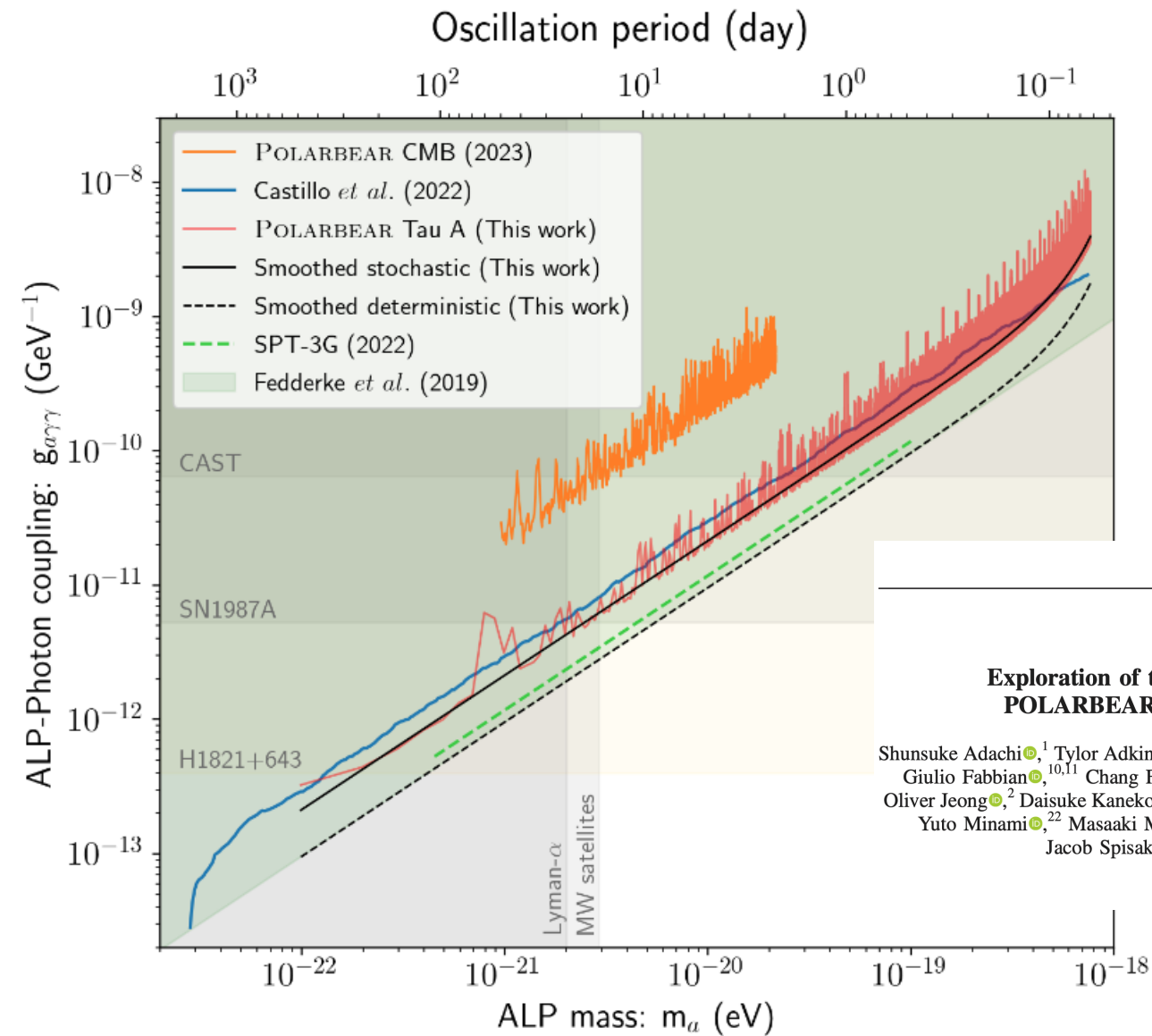
POLARBEAR 2023

- Test for presence of signal:

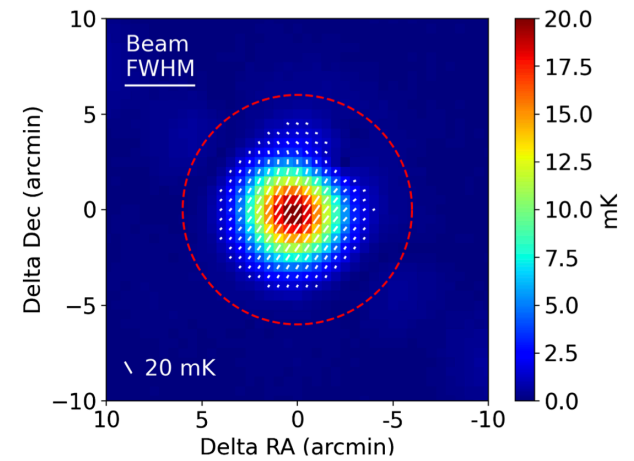
$$\Delta\chi^2 \equiv \chi^2(A=0) - \chi^2(A^{\text{mle}}, f^{\text{mle}}, \theta^{\text{mle}})$$

- Compare to a simulated distribution
- $\sigma_{\text{PTE}} = 1.7$ : no significant detection
- Place 95% upper confidence limit on sinusoid amplitude  $A_{95}$  across frequency range





Adachi et al Crab Nebula 2024



PHYSICAL REVIEW D **110**, 063013 (2024)

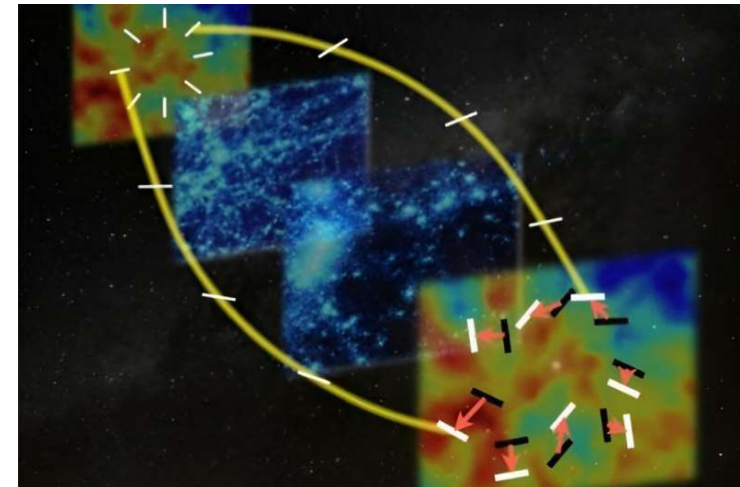
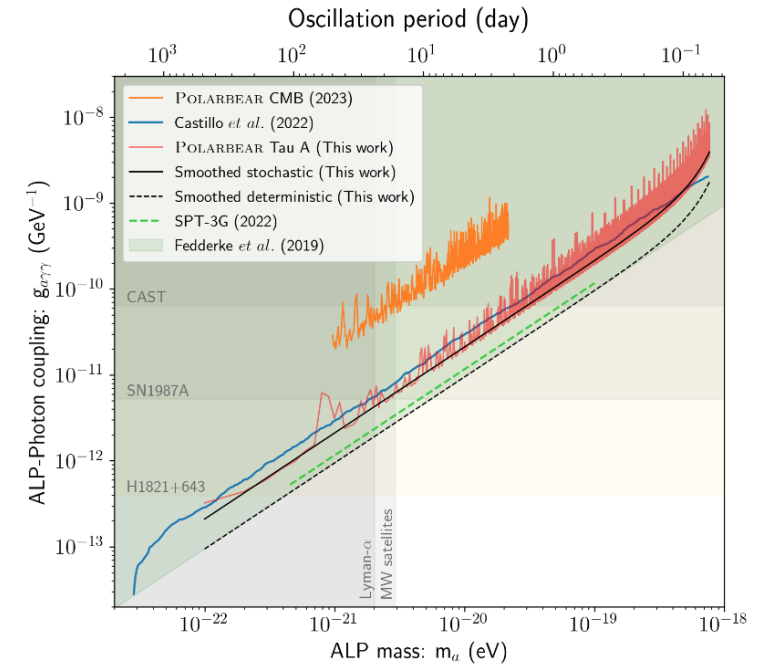
### Exploration of the polarization angle variability of the Crab Nebula with POLARBEAR and its application to the search for axionlike particles

Shunsuke Adachi<sup>1</sup>, Tylor Adkins<sup>2</sup>, Carlo Baccigalupi<sup>3,4,5</sup>, Yuji Chinone<sup>6,7</sup>, Kevin T. Crowley<sup>8</sup>, Josquin Errard<sup>9</sup>, Giulio Fabbian<sup>10,11</sup>, Chang Feng<sup>12,13</sup>, Takuro Fujino<sup>14</sup>, Masaya Hasegawa<sup>15,6,16</sup>, Masashi Hazumi<sup>6,15,7,17,16</sup>, Oliver Jeong<sup>2</sup>, Daisuke Kaneko<sup>6</sup>, Brian Keating<sup>8</sup>, Akito Kusaka<sup>18,19,7,20</sup>, Adrian T. Lee<sup>2,19</sup>, Anto I. Lonappan<sup>21</sup>, Yuto Minami<sup>22</sup>, Masaaki Murata<sup>18</sup>, Lucio Piccirillo<sup>23</sup>, Christian L. Reichardt<sup>24</sup>, Praween Siritanasak<sup>25</sup>, Jacob Spisak<sup>8</sup>, Satoru Takakura<sup>26</sup>, Grant P. Teply<sup>8</sup>, and Kyohei Yamada<sup>18,\*</sup>

(POLARBEAR Collaboration)

# Conclusion

- Well-motivated from fundamental physics and from dark matter
- Constraints coming in
- Reminds me of early days of BICEP and inflation searches
- First they ignore you, then they laugh at you, then they fight you, then they join you, then you win.



# END



Data 🖱️

<https://briankeating.com/cb>

